

# Math 562 Fall 2020

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# Outline

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- 1 **General Info**
- 2 5.5 Girsanov's Theorem
- 3 5.6 A Few Applications of Girsanov's Theorem

HW4 has been graded.

I will set up HW5 later today. The due date is Nov. 6 at noon.

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### Theorem 5.22 (Girsanov's theorem)

Assume that the probability measures  $\mathbb{P}$  and  $\mathbb{Q}$  are mutually absolutely continuous on  $\mathcal{F}_\infty$ . Let  $(D_t)_{t \geq 0}$  be the martingale with cadlag sample paths such that, for every  $t \geq 0$ ,

$$D_t = \frac{d\mathbb{Q}}{d\mathbb{P}} \Big|_{\mathcal{F}_t}.$$

Assume that  $D$  has continuous sample paths, and let  $L$  be the unique continuous local martingale such that  $D_t = \mathcal{E}(L)_t$ . Then, if  $M$  is a continuous local martingale under  $\mathbb{P}$ , the process

$$\tilde{M} = M - \langle M, L \rangle$$

is a continuous local martingale under  $\mathbb{Q}$ .

In most applications of Girsanov's theorem, one constructs the probability measure  $\mathbb{Q}$  in the following way. Start from a continuous local martingale  $L$  such that  $L_0 = 0$  and  $\langle L, L \rangle_\infty < \infty$  a.s. The latter condition implies that the limit  $L_\infty := \lim_{t \rightarrow \infty} L_t$  exists a.s. Then  $\mathcal{E}(L)_t$  is a nonnegative continuous local martingale hence a supermartingale, which converges a.s. to

$$\mathcal{E}(L)_\infty = \exp\left(L_\infty - \frac{1}{2}\langle L, L \rangle_\infty\right),$$

and  $\mathbb{E}[\mathcal{E}(L)_\infty] \leq 1$  by Fatou's lemma. If the property

$$\mathbb{E}[\mathcal{E}(L)_\infty] = 1 \tag{1}$$

holds, then  $\mathcal{E}(L)$  is a uniformly integrable martingale (by Fatou's lemma again, one has  $\mathcal{E}(L)_t \geq \mathbb{E}[\mathcal{E}(L)_\infty | \mathcal{F}_t]$ , but (1) implies that  $\mathbb{E}[\mathcal{E}(L)_\infty] = \mathbb{E}[\mathcal{E}(L)_0] = \mathbb{E}[\mathcal{E}(L)_t]$  for every  $t \geq 0$ ). If we let  $\mathbb{Q}$  be the probability measure with density  $\mathcal{E}(L)_\infty$ , with respect to  $\mathbb{P}$ , we are in the setting of Theorem 5.22, with  $D_t = \mathcal{E}(L)_t$ . It is therefore very important to give conditions that ensure that (1) holds.



### Theorem 5.23

Let  $L$  be a continuous local martingale such that  $L_0 = 0$ . Consider the following properties:

- (i)  $\mathbb{E}[\exp(\frac{1}{2}\langle L, L \rangle_\infty)] < \infty$  (Novikov's condition);
- (ii)  $L$  is a uniformly integrable martingale, and  $\mathbb{E}[\exp(\frac{1}{2}L_\infty)] < \infty$  (Kazamaki's condition);
- (iii)  $\mathcal{E}(L)$  is a uniformly integrable martingale.

Then (i)  $\Rightarrow$  (ii)  $\Rightarrow$  (iii).

### Proof of Theorem 5.23

(i) $\Rightarrow$ (ii): Property (i) implies that  $\mathbb{E}[\langle L, L \rangle_\infty] < \infty$  hence also that  $L$  is a continuous martingale bounded in  $L^2$ . Then,

$$\exp\left(\frac{1}{2}L_\infty\right) = (\mathcal{E}(L)_\infty)^{1/2}(\exp\left(\frac{1}{2}\langle L, L \rangle_\infty\right))^{1/2}$$

so that, by the Cauchy-Schwarz inequality,

$$\begin{aligned}\mathbb{E}[\exp\left(\frac{1}{2}L_\infty\right)] &\leq (\mathbb{E}[\mathcal{E}(L)_\infty])^{1/2}(\mathbb{E}[\exp\left(\frac{1}{2}\langle L, L \rangle_\infty\right)])^{1/2} \\ &\leq (\mathbb{E}[\exp\left(\frac{1}{2}\langle L, L \rangle_\infty\right)])^{1/2} < \infty.\end{aligned}$$

### Proof of Theorem 5.23 (cont)

(ii) $\Rightarrow$ (iii): Since  $L$  is a uniformly integrable martingale, Theorem 3.22 shows that, for any stopping time  $T$ , we have  $L_T = \mathbb{E}[L_\infty | \mathcal{F}_T]$ . Jensen's inequality then gives

$$\exp\left(\frac{1}{2}L_T\right) \leq \mathbb{E}\left[\exp\left(\frac{1}{2}L_\infty\right) | \mathcal{F}_T\right].$$

By assumption,  $\mathbb{E}[\exp(\frac{1}{2}L_\infty)] < \infty$  which implies that the collection of all variables of the form  $\mathbb{E}[\exp(\frac{1}{2}L_\infty) | \mathcal{F}_T]$ , for any stopping time  $T$ , is uniformly integrable. The preceding bound then shows that the collection of all variables  $\exp(\frac{1}{2}L_T)$ , for any stopping time  $T$ , is also uniformly integrable.

For  $0 < a < 1$ , define  $Z_t^{(a)} = \exp(\frac{aL_t}{1+a})$ . Then, one easily verifies that

$$\mathcal{E}(aL)_t = (\mathcal{E}(L)_t)^{a^2} (Z_t^{(a)})^{1-a^2}.$$

### Proof of Theorem 5.23 (cont)

If  $\Gamma \in \mathcal{F}$  and  $T$  is a stopping time, Hölder's inequality gives

$$\begin{aligned}\mathbb{E}[1_{\Gamma} \mathcal{E}(aL)_T] &\leq (\mathbb{E}[\mathcal{E}(L)_T])^{a^2} (\mathbb{E}[1_{\Gamma} Z_T^{(a)}])^{1-a^2} \\ &\leq (\mathbb{E}[1_{\Gamma} Z_T^{(a)}])^{1-a^2} \leq (\mathbb{E}[1_{\Gamma} \exp(\frac{1}{2}L_T)])^{2a(1-a)}.\end{aligned}$$

In the second inequality, we used the property  $\mathbb{E}[\mathcal{E}(L)_T] \leq 1$ , which holds by Proposition 3.25 because  $\mathcal{E}(L)$  is a nonnegative supermartingale and  $\mathcal{E}(L)_0 = 1$ . In the third inequality, we use Jensen's inequality, noting that  $\frac{1+a}{2a} > 1$  since the collection of all variables of the form  $\exp(\frac{1}{2}L_T)$ , for any stopping time  $T$ , is uniformly integrable, the preceding display shows that so is the collection of all variables  $\mathcal{E}(aL)_T$ , for any stopping time  $T$ .

### Proof of Theorem 5.23 (cont)

By the definition of a continuous local martingale, there is an increasing sequence  $T_n \uparrow \infty$  of stopping times, such that, for every  $n$ ,  $\mathcal{E}(aL)_{t \wedge T_n}$  is a martingale. If  $0 \leq s \leq t$ , we can use uniform integrability to pass to the limit  $n \rightarrow \infty$  in the equality  $\mathbb{E}[\mathcal{E}(aL)_{t \wedge T_n} | \mathcal{F}_s] = \mathcal{E}(aL)_{s \wedge T_n}$  and we get that  $\mathcal{E}(aL)$  is a uniformly integrable martingale. It follows that

$$\begin{aligned} 1 &= \mathbb{E}[\mathcal{E}(aL)_\infty] \leq (\mathbb{E}[\mathcal{E}(L)_\infty])^{a^2} (\mathbb{E}[Z_\infty^{(a)}])^{1-a^2} \\ &\leq (\mathbb{E}[\mathcal{E}(L)_\infty])^{a^2} (\mathbb{E}[\exp(\frac{1}{2}L)_\infty])^{2a(1-a)}, \end{aligned}$$

using again Jensen's inequality as above. When  $a \uparrow 1$ , this gives  $\mathbb{E}[\mathcal{E}(L)_\infty] \geq 1$  hence  $\mathbb{E}[\mathcal{E}(L)_\infty] = 1$ .

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In this section, we describe a few applications of Girsanov's theorem, which illustrate the strength of the previous results.

### Constructing solutions of stochastic differential equations

Let  $b$  be a bounded measurable function on  $\mathbb{R}_+ \times \mathbb{R}$ . We assume that there exists a function  $g \in L^2(\mathbb{R}_+, \mathcal{B}(\mathbb{R}_+), dt)$  such that  $|b(t, x)| \leq g(t)$  for every  $(t, x) \in \mathbb{R}_+ \times \mathbb{R}$ . This holds in particular if there exists an  $A > 0$  such that  $|b|$  is bounded on  $[0, A] \times \mathbb{R}$  and vanishes on  $(A, \infty) \times \mathbb{R}$ .

Let  $B$  be an  $(\mathcal{F}_t)$ -Brownian motion. Consider the continuous local martingale

$$L_t = \int_0^t b(s, B_s) dB_s$$

and the associated exponential martingale

$$D_t = \mathcal{E}(L)_t = \exp \left( \int_0^t b(s, B_s) dB_s - \frac{1}{2} \int_0^t b^2(s, B_s) ds \right).$$

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Our assumption on  $b$  ensures that condition (i) of Theorem 5.23 holds, and thus  $D$  is a uniformly integrable martingale. We set  $\mathbb{Q} = D_\infty \cdot \mathbb{P}$ . Girsanov's theorem, and remark (c) following the statement of the theorem, show that the process

$$\beta_t := B_t - \int_0^t b(s, B_s) ds$$

is an  $(\mathcal{F}_t)$ -Brownian motion under  $\mathbb{Q}$ .

We can restate the latter property by saying that, under the probability measure  $\mathbb{Q}$ , there exists an  $(\mathcal{F}_t)$ -Brownian motion  $\beta$  such that the process  $X = B$  solves the stochastic differential equation

$$dX_t = d\beta_t + b(t, B_t)dt.$$

This equation is of the type that will be considered in Chap. 7 below. Note that we are not making any regularity assumption on the function  $b$ . It is remarkable that Girsanov's theorem allows one to construct solutions of stochastic differential equations without regularity conditions on the coefficients.

## The Cameron-Martin formula

We now assume that  $b(t, x) = g(t)$ , where  $g \in L^2(\mathbb{R}_+, \mathcal{B}(\mathbb{R}_+), dt)$ , and we also set, for every  $t \geq 0$ ,

$$h(t) = \int_0^t g(s) ds.$$

The set  $\mathcal{H}$  of all functions  $h$  that can be written in this form is called the Cameron-Martin space. If  $h \in \mathcal{H}$ , we sometimes write  $\dot{h} = g$  for the associated function in  $L^2(\mathbb{R}_+, \mathcal{B}(\mathbb{R}_+), dt)$  (this is the derivative of  $h$  in the sense of distributions).

As a special case of the previous discussion, under the probability measure

$$\mathbb{Q} := D_\infty \cdot \mathbb{P} = \exp\left(\int_0^\infty g(t) dB_t - \frac{1}{2} \int_0^\infty g^2(t) dt\right) \cdot \mathbb{P},$$

the process  $\beta_t = B_t - h(t)$  is a Brownian motion. Hence, for every

non-negative measurable function  $\Phi$  on  $C(\mathbb{R}_+, \mathbb{R})$ ,

$$\begin{aligned}\mathbb{E}_{\mathbb{P}}[D_{\infty} \cdot \Phi((B_t)_{t \geq 0})] &= \mathbb{E}_{\mathbb{Q}}[\Phi((B_t)_{t \geq 0})] \\ &= \mathbb{E}_{\mathbb{Q}}[\Phi((\beta_t + h(t))_{t \geq 0})] = \mathbb{E}_{\mathbb{P}}[\Phi((B_t + h(t))_{t \geq 0})].\end{aligned}$$

The equality between the two ends of the last display is the Cameron-Martin formula. In the next proposition, we write this formula in the special case of the canonical construction of Brownian motion on the Wiener space.

### Proposition 5.24 (Cameron-Martin formula)

Let  $W(dw)$  be the Wiener measure on  $C(\mathbb{R}_+, \mathbb{R})$ , and let  $h$  be a function in the Cameron-Martin space  $\mathcal{H}$ . Then, for every non-negative measurable function  $\Phi$  on  $C(\mathbb{R}_+, \mathbb{R})$ ,

$$\int W(dw) \Phi(w+h) = \int W(dw) \exp\left(\int_0^{\infty} h(s) dw_s - \frac{1}{2} \int_0^{\infty} h(s)^2 ds\right) \Phi(w).$$

non-negative measurable function  $\Phi$  on  $C(\mathbb{R}_+, \mathbb{R})$ ,

$$\begin{aligned}\mathbb{E}_{\mathbb{P}}[D_{\infty} \cdot \Phi((B_t)_{t \geq 0})] &= \mathbb{E}_{\mathbb{Q}}[\Phi((B_t)_{t \geq 0})] \\ &= \mathbb{E}_{\mathbb{Q}}[\Phi((\beta_t + h(t))_{t \geq 0})] = \mathbb{E}_{\mathbb{P}}[\Phi((B_t + h(t))_{t \geq 0})].\end{aligned}$$

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## Remark

The integral  $\int_0^\infty \dot{h}(s)dw_s$  is a stochastic integral with respect to  $w(s)$  (which is a Brownian motion under  $W(dw)$ ), but it can also be viewed as a Wiener integral since the function  $\dot{h}$  is deterministic. The Cameron-Martin formula can be established by Gaussian calculations that do not involve stochastic integrals or Girsanov's theorem. Still it is instructive to derive this formula as a special case of Girsanov's theorem.

The Cameron-Martin formula gives a “quasi-invariance” property of Wiener measure under the translations by functions of the Cameron-Martin space: The image of the Wiener measure  $W(dw)$  under the mapping  $w \mapsto w + h$  has a density with respect to  $W(dw)$  and this density is the terminal value of the exponential martingale associated with the martingale  $\int_0^t \dot{h}(s)dw_s$ .

## Law of hitting times for Brownian motion with drift

Let  $B$  be a 1-dim Brownian motion with  $B_0 = 0$ , and for every  $a > 0$ , let  $T_a := \inf\{t \geq 0 : B_t = a\}$ . If  $c \in \mathbb{R}$ , we would like to compute the law of the stopping time

$$U_a := \inf\{t \geq 0 : B_t + ct = a\}.$$

Of course, if  $c = 0$ , we have  $U_a = T_a$ , and the desired distribution is given by Corollary 2.22. Girsanov's theorem (or rather the Cameron-Martin formula) will allow us to derive the case where  $c$  is arbitrary from the special case  $c = 0$ .

Fix  $t > 0$  and apply the Cameron-Martin formula with

$$h(s) = c1_{\{s \leq t\}}, \quad h(s) = c(s \wedge t).$$

and, for every  $w \in C(\mathbb{R}_+, \mathbb{R})$ ,

$$\Phi(w) = 1_{\{\max_{[0,t]} w(s) > a\}}.$$

It follows that

$$\begin{aligned}
 \mathbb{P}(U_a \leq t) &= \mathbb{E}[\Phi(B + h)] \\
 &= \mathbb{E}\left[\Phi(B) \exp\left(\int_0^\infty h(s)dB_s - \frac{1}{2} \int_0^\infty h(s)^2 ds\right)\right] \\
 &= \mathbb{E}[1_{\{T_a \leq t\}} \exp(cB_t - \frac{c^2}{2}t)] = \mathbb{E}[1_{\{T_a \leq t\}} \exp(cB_{t \wedge T_a} - \frac{c^2}{2}(t \wedge T_a))] \\
 &= \mathbb{E}[1_{\{T_a \leq t\}} \exp(ca - \frac{c^2}{2}T_a)] = \int_0^t ds \frac{a}{\sqrt{2\pi s^3}} e^{-\frac{a^2}{2s}} e^{ca - \frac{c^2}{2}s} \\
 &= \int_0^t ds \frac{a}{\sqrt{2\pi s^3}} e^{-\frac{1}{2s}(a - cs)^2}
 \end{aligned}$$

where, in the fourth equality, we used the optional stopping theorem to write

$$\mathbb{E}[\exp(cB_t - \frac{c^2}{2}t) | \mathcal{F}_{t \wedge T_a}] = \exp(cB_{t \wedge T_a} - \frac{c^2}{2}(t \wedge T_a))$$

and we also made use of the explicit density of  $T_a$  given in Corollary 2.22. This calculation shows that the variable  $U_a$  has a density on  $\mathbb{R}_+$  given by

$$\psi(s) = \frac{a}{\sqrt{2\pi s^3}} e^{-\frac{1}{2s}(a-cs)^2}.$$

By integrating this density, we get that

$$\mathbb{P}(U_a < \infty) = \begin{cases} 1, & \text{if } c \geq 0; \\ e^{2ca}, & \text{if } c \leq 0, \end{cases}$$

which may also be checked more easily by applying the optional stopping theorem to the continuous martingale  $\exp(-2c(B_t + ct))$ .